

Measurement of III–V Compound Semiconductor Characteristics using the Contactless Electroreflectance Method

Jae-In Yu*, Soon-Don Choi[†] and Ho-Gyeong Chang**

Abstract – The electromodulation methods of photoreflectance and the related technique of contactless electroreflectance (CER) are valuable tools in the evaluation of important device parameters for structures such as heterojunction bipolar transistors, pseudomorphic high electron mobility transistors, and quantum dots (QDs). CER is a very general principle of experimental physics. Instead of measuring the optical reflectance of the material, the derivative with respect to a modulating electric field is evaluated. This procedure generates sharp, differential-like spectra in the region of interband (intersubband) transitions. We conduct electric-optical studies of both GaAs layers and InAs self-assembled QDs grown by molecular beam epitaxy. Strong GaAs bandgap energy is measured in both structures. In the case of InAs monolayers in GaAs matrices, the strong GaAs bandgap energy is caused by the lateral quantum confinement.

Keywords: InAs, QD, CER, GaAs, Semiconductor

1. Introduction

Semiconductor quantum dots (QDs) have attracted considerable interest recently due to their potential in technological applications. Although Stranski–Krastanow growth is regarded as a promising road toward zero-dimensional quantum structures, its random nucleation in a single layer results in a broad distribution of QD size and position. QD arrays of the same size and shape are required for practical applications; thus, the identification of growth mechanisms that lead to a narrowing of the size distribution is of great importance [1, 2]. Modulation reflectance spectroscopy is a powerful tool that investigates the optical properties of III–V compound semiconductors. Since the modulation reflectivity is an absorption-type technique, the acquired information reflects the optical transition behavior, which is more directly related to the theoretical absorption model. The optical properties of InAs/GaAs QDs are studied by contactless electroreflectance (CER) and scanning electron microscopy.

2. Theory

Modulation spectroscopy is an important technique for the study and characterization of semiconductor thin films and heterostructures. Techniques such as CER and

photoreflectance are very useful in probing the bandstructure of the semiconductor because they measure the differential reflectance change and are sensitive to surface or interface electric fields. They are particularly useful techniques for probing the band structure of the semiconductors when conventional techniques such as photoluminescence cannot be used because of its bad sample quality. CER requires no special sample mounting and can be performed in a variety of transparent ambience because it is contactless. In CER measurement, the sample is placed between two capacitor plates. The size of the spacer is such that there is a very thin layer of air between the front surface of the sample and the conducting part of the first electrode. There is nothing in direct contact with the front surface of the sample. The CER spectra as a function of photon energy can be fitted using a familiar Aspnes derivative function in the low electric field limit [3–5]

$$\frac{\Delta R}{R} = \text{Re} \sum_{j=1}^p C_j e^{i\theta_j} (E - E_{g_j} + i\Gamma_j)^{-n}$$

In the above equation, R is the reflectance, ΔR is the induced change in the reflectance by modulation light, E is the photon energy, and p is the total number of spectral structures to be fitted. E_{g_j} , Γ_j , C_j , and θ_j are the transition energy, broadening parameter, amplitude, and phase of the j -th feature corresponding to a critical point, respectively.

3. Experiment

The InAs/GaAs QDs structure was grown on undoped GaAs substrates by molecular beam epitaxy system. After the growth

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of a 500nm-thick GaAs buffer layer at the substrate temperature (T_g) of 540°C, InAs/GaAs QDs structures were grown at T_g of 460°C.

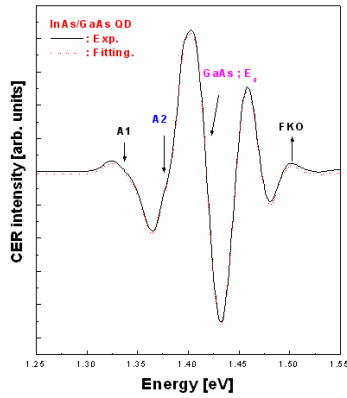
The growth sequence for the InAs/GaAs QDs structure is presented in Table 1.

Table 1. The growth structure of InAs/GaAs QDs.

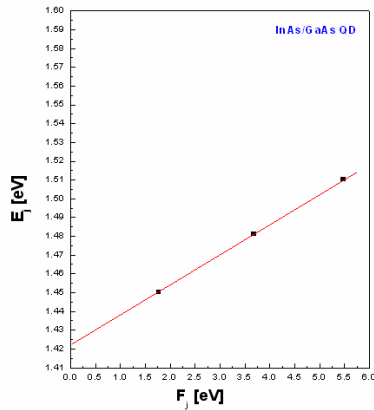
GaAs 90 nm	0.5 ML/s
GaAs 10 nm	0.5 ML/s
As supply	1×10^{-4} Torr
In 4.5 ML	0.142 ML/s
Ga 1 ML	0.1 ML/s
GaAs buffer 500 nm	0.8 ML/s
UndopedGaAs sub	

4. Result and Discussion

We begin by analyzing the morphology of the WL formed on deposition of 4.5 ML of InAs on the GaAs(001) surface. Fig. 1 shows the transitions of A1(1.34 eV) and A2(1.37 eV) in CER spectra of the InAs/GaAs QD sample, which are ascribed to the heavy-hole-type and the light-hole-type transitions, respectively.



(a)



(b)

Fig. 1. The illustrated CER spectrum in InAs/GaAs QDs.

We have

$$\left[\frac{1}{r^2} \left(r \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{\partial^2}{\partial \theta^2} \right) - \frac{\partial^2}{\partial z^2} + V_e(r, z) \right] \Psi(r, \theta, z) = E \Psi(r, \theta, z)$$

Wave function can be given by two sub-functions, slowly varying function $g_r(z)$ of r and $f_m(r)$, which fit a set of equations for each angular momentum channel m :

$$\begin{aligned} \left[-\frac{\partial^2}{\partial z^2} + V_e(r, z) \right] g_r &= E_o(R) g_r(z), \\ \left[-\frac{1}{r^2} \left(r \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + m^2 \right) - E_o(r) \right] f_m(r) &= E f_m(r) \end{aligned}$$

Then, the potential distribution of the conductive band-edge within the eight-band KP theory model derived by Califano et al. [6, 7] can be expressed as

$$V_e = E_{v,av} + \frac{\Delta_o}{3} + E_{g,o} + a_e e_{hh}$$

In addition, that of the heavy hole is

$$V_{hh} = E_{v,av} + \frac{\Delta_o}{3} + a_e e_{hh} - b e_t$$

where $E_{v,av}$ is the average valence-band energy; Δ_o is the spin-orbit splitting; a_e and b are the deformation potentials; and e_{hh} and e_t are the hydrostatic and uniaxial strain, respectively. This result is shown in Table 2.

Table 2. Fitting experimental values and calculation results with effective mass approximation.

InAs/GaAs QDs		
Peak	CER [Transition E(eV);300K]	Calculation [Transition (eV);300K]
A1	1.344	1.347
A2	1.375	1.382

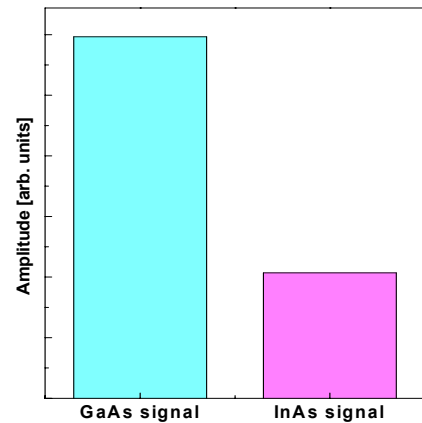


Fig. 2. The graph of amplitude in GaAs and InAs signals

The strong GaAs bandgap energy was measured. In the case of the InAs monolayer in GaAs matrices, the strong GaAs bandgap energy is caused by the lateral quantum confinement. For the InAs QD sample, additional spectral features appeared as shown in Fig. 3. The critical point (CP) in the CER spectrum at 1.405 eV exhibits an emission of red-shift of ~ 11 meV compared with that of the annealing InAs QD sample. The annealing sample InAs QD disappeared at the A1 and A2 peaks. This result was caused by surface damage [8–12].

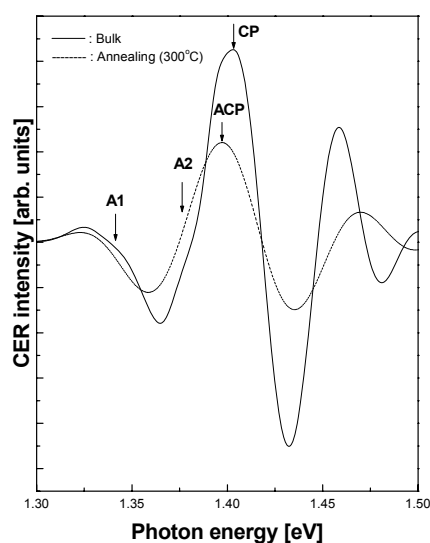


Fig. 3. The CER spectra in InAs QD and annealing InAs QD sample.

Fig. 4 shows an optical photograph of the annealing InAs QD sample (100–400 °C).

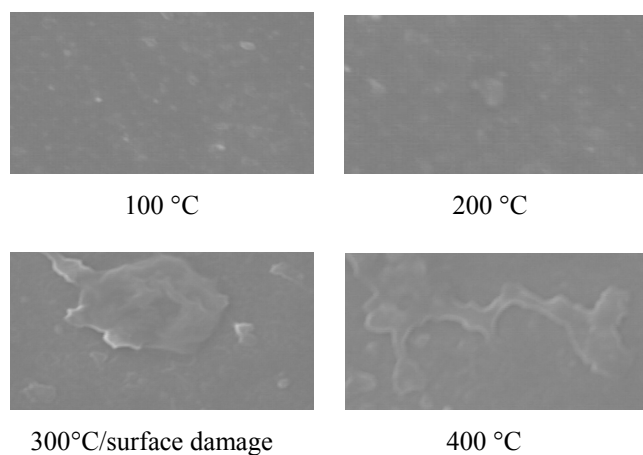


Fig. 4. The photograph of the annealing InAs QD sample.

4. Conclusion

We investigated the CER spectra in the InAs QD sample. The CER result and the strong GaAs bandgap energy were measured. In the case of the InAs monolayer in GaAs matrices, the strong GaAs bandgap energy was caused by the lateral quantum confinement. The CP in the CER spectrum at 1.405 eV exhibited an emission of red-shift of ~ 11 meV compared with that of the annealing InAs QD sample. This result was caused by surface damage.

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