# A CERTAIN KÄHLER POTENTIAL OF THE POINCARÉ METRIC AND ITS CHARACTERIZATION 

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#### Abstract

We will show a rigidity of a Kähler potential of the Poincaré metric with a constant length differential.


## 1. Introduction

From the fundamental result of Donnelly-Fefferman [4], the vanishing of the space of $L^{2}$ harmonic $(p, q)$ forms has been an important research theme in the theory of complex domains. Since M. Gromov ([6], see also [2]) suggested the concept of the Kähler hyperbolicity and gave a connection to the vanishing theorem, there have been many studies on the Kähler hyperbolicity of the Bergman metric, which is a fundamental Kähler structure of bounded pseudoconvex domains. The Kähler structure $\omega$ is Kähler hyperbolic if there is a global 1-form $\eta$ with $d \eta=\omega$ and $\sup \|\eta\|_{\omega}<\infty$.

In [3], H. Donnelly showed the Kähler hyperbolicity of Bergman metric on some class of weakly pseudoconvex domains. For bounded homogeneous domain $D$ in $\mathbb{C}^{n}$ and its Bergman metric $\omega_{D}$ especially, he used a classical result of Gindikin [5] to show that sup $\left\|d \log K_{D}\right\|_{\omega_{D}}<\infty$. Here $K_{D}$ is the Bergman kernel function of $D$ so $\log K_{D}$ is a canonical potential of $\omega_{D}$.

In their paper [7], S. Kai and T. Ohsawa gave another approach. They proved that every bounded homogeneous domain has a Kähler potential of the Bergman metric whose differential has a constant length.

Theorem 1.1 (Kai-Ohsawa [7]). For a bounded homogeneous domain $D$ in $\mathbb{C}^{n}$, there exists a positive real valued function $\varphi$ on $D$ such that $\log \varphi$ is a Kähler potential of the Bergman metric $\omega_{D}$ and $\|d \log \varphi\|_{\omega_{D}}$ is constant.

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It can be obtained by the facts that each homogeneous domain is biholomorphic to a Siegel domain (see [10]) and a homogeneous Siegel domain is affine homogeneous (see [8]).

More precisely, let us consider a bounded homogeneous domain $D$ in $\mathbb{C}^{n}$ and a biholomorphism $F: D \rightarrow S$ for a Siegel domain $S$. For the Bergman kernel function $K_{S}$ of $S$ which is a canonical potential of the Bergman metric $\omega_{S}$, it is easy to show that $d \log K_{S}$ has a constant length with respect to $\omega_{S}$ from the affine homogeneity of $S$ (the group of affine holomorphic automorphisms acts transitively on $S$ ). Since $\log K_{S}$ is a Kähler potential of $\omega_{S}$, the transformation formula of the Bergman kernel implies that the pullback $F^{*} \log K_{S}=\log K_{S} \circ F$ is also a Kähler potential of $\omega_{D}$. Using the fact that $F:\left(D, \omega_{D}\right) \rightarrow\left(S, \omega_{S}\right)$ is an isometry, we have $\left\|d\left(F^{*} \log K_{S}\right)\right\|_{\omega_{D}}=\left\|d \log K_{S}\right\|_{\omega_{S}} \circ F$. As a function $\varphi$ in Theorem 1.1, we can choose the pullback $K_{S} \circ F$ of the Bergman kernel function of the Siegel domain.

At this junction, it is natural to ask:
If there is a Kähler potential $\log \varphi$ with a constant $\|d \log \varphi\|_{\omega_{D}}$, is it always obtained by the pullback of the Bergman kernel function of the Siegel domain?
The aim of this paper is to discuss of this question in the 1-dimensional case.
The only bounded homogeneous domain in $\mathbb{C}$ is the unit disc $\Delta=\{z \in \mathbb{C}$ : $|z|<1\}$ up to the biholomorphic equivalence and the 1-dimensional correspondence of the Bergman metric, namely a holomorphically invariant hermitian structure, is only the Poincaré metric. Hence the main theorem as follows gives a positive answer to the question.

Theorem 1.2. Let $\omega_{\Delta}$ be the Poincaré metric of the unit disc $\Delta$. Suppose that there exists a positive real valued function $\varphi: \Delta \rightarrow \mathbb{R}$ such that $\log \varphi$ is a Kähler potential of the Poincaré metric and $\|d \log \varphi\|_{\omega_{\Delta}}$ is constant on $\Delta$. Then $\varphi$ is the pullback of the canonical potential on the half-plane $\mathbf{H}=\{z \in \mathbb{C}: \operatorname{Re} z<0\}$.

Note that 1-dimensional Siegel domain is just the half-plane. We will introduce the Poincaré metric and related notions in Section 2. As an application of the main theorem, we can characterize the half-plane by the canonical potential.

Corollary 1.3. Let $D$ be a simply connected, proper domain in $\mathbb{C}$ with a Poincaré metric $\omega_{D}=i \lambda d z \wedge d \bar{z}$. If $\|d \log \lambda\|_{\omega_{D}}$ is constant on $D$, then $D$ is affine equivalent to the half-plane $\mathbf{H}=\{z \in \mathbb{C}: \operatorname{Re} z<0\}$.

In Section 2, we will introduce notions and concrete version of the main theorem. Then we will study the existence of a nowhere vanishing complete holomorphic vector field which is tangent to a potential whose differential is of constant length (Section 3). Using relations between complete holomorphic vector fields and model potentials in Section 4, we will prove theorems.

## 2. Background materials

Let $X$ be a Riemann surface. The Poincaré metric of $X$ is a complete hermitian metric with a constant Gaussian curvature, -4 . The Poincaré metric exists on $X$ if and only if $X$ is a quotient of the unit disc. If $X$ is covered by $\Delta$, the Poincaré metric can be induced by the covering map $\pi: \Delta \rightarrow X$ and it is uniquely determined. Throughout of this paper, the Kähler form of the Poincaré metric of $X$, denoted by $\omega_{X}$, stands for the metric also. When $\omega_{X}=i \lambda d z \wedge d \bar{z}$ in the local holomorphic coordinate function $z$, the curvature can be written by

$$
\kappa=-\frac{2}{\lambda} \frac{\partial^{2}}{\partial z \partial \bar{z}} \log \lambda
$$

So the curvature condition $\kappa \equiv-4$ implies that

$$
\frac{\partial^{2}}{\partial z \partial \bar{z}} \log \lambda=2 \lambda
$$

equivalently

$$
d d^{c} \log \lambda=2 \omega_{X}
$$

where $d^{c}=\frac{i}{2}(\bar{\partial}-\partial)$. That means the function $\frac{1}{2} \log \lambda$ is a local Kähler potential of $\omega_{X}$. Any other local potential of $\omega_{X}$ is always of the form $\frac{1}{2} \log \lambda+\log |f|^{2}$ where $f$ is a local holomorphic function on the domain of $z$. We call $\frac{1}{2} \log \lambda$ the canonical potential with respect to the coordinate function $z$. For a domain $D$ in $\mathbb{C}$, the canonical potential of $D$ means the canonical potential with respect to the standard coordinate function of $\mathbb{C}$.

Let us consider the Poincaré metric $\omega_{\Delta}$ of the unit disc $\Delta$ :

$$
\omega_{\Delta}=i \frac{1}{\left(1-|z|^{2}\right)^{2}} d z \wedge d \bar{z}=i \lambda_{\Delta} d z \wedge d \bar{z}
$$

The canonical potential $\lambda_{\Delta}$ satisfies

$$
\begin{aligned}
\left\|d \log \lambda_{\Delta}\right\|_{\omega_{\Delta}}^{2} & =\left\|\frac{\partial \log \lambda_{\Delta}}{\partial z} d z+\frac{\partial \log \lambda_{\Delta}}{\partial \bar{z}} d \bar{z}\right\|_{\omega_{\Delta}}^{2} \\
& =\frac{\partial \log \lambda_{\Delta}}{\partial z} \frac{\partial \log \lambda_{\Delta}}{\partial \bar{z}} \frac{1}{\lambda_{\Delta}}=4|z|^{2}
\end{aligned}
$$

so does not have a constant length. By the same way of Kai-Ohsawa [7], we can get a model for $\varphi$ in Theorem 1.1 for the unit disc,

$$
\begin{equation*}
\varphi_{\theta}(z)=\frac{\left|1+e^{i \theta} z\right|^{4}}{\left(1-|z|^{2}\right)^{2}} \quad \text { for } \theta \in \mathbb{R} \tag{2.1}
\end{equation*}
$$

as a pullback of the canonical potential $\lambda_{\mathbf{H}}=1 /|\operatorname{Re} w|^{2}$ on the left-half plane $\mathbf{H}=\{w: \operatorname{Re} w<0\}$ by the Cayley transforms (see (4.3) for instance). The
term $\theta$ depends on the choice of the Cayley transform. Since $\log \varphi_{\theta}=\log \lambda_{\Delta}+$ $\log \left|1+e^{i \theta} z\right|^{4}$, the function $\frac{1}{2} \log \varphi_{\theta}$ is a Kähler potential. Moreover

$$
\left\|d \log \varphi_{\theta}\right\|_{\omega_{\Delta}}^{2} \equiv 4
$$

At this moment, we introduce a significant result of Kai-Ohsawa.
Theorem 2.1 (Kai-Ohsawa [7]). For a bounded homogeneous domain $D$ in $\mathbb{C}^{n}$, suppose that there is a Kähler potential $\log \psi$ of the Bergman metric $\omega_{D}$ with a constant $\|d \log \psi\|_{\omega_{D}}$, then $\|d \log \psi\|_{\omega_{D}}=\|d \log \varphi\|_{\omega_{D}}$ where $\varphi$ is as in Theorem 1.1.

Suppose that a positively real valued $\varphi$ on $\Delta$ satisfies that $d d^{c} \log \varphi=2 \omega_{\Delta}$ and $\|d \log \varphi\|_{\omega_{\Delta}}^{2} \equiv c$ for some constant $c$. Theorem 2.1 implies that $c$ must be 4. Therefore, we can rewrite Theorem 1.2 by:

Theorem 2.2. If there exists a function $\varphi: \Delta \rightarrow \mathbb{R}$ satisfying

$$
\begin{equation*}
d d^{c} \log \varphi=2 \omega_{\Delta} \quad \text { and } \quad\|d \log \varphi\|_{\omega_{\Delta}}^{2} \equiv 4 \tag{2.2}
\end{equation*}
$$

then $\varphi=r \varphi_{\theta}$ as in (2.1) for some $r>0$ and $\theta \in \mathbb{R}$.
Corollary 1.3 can be also written by:
Corollary 2.3. Let $D$ be a simply connected, proper domain in $\mathbb{C}$ with a Poincaré metric $\omega_{D}=i \lambda d z \wedge d \bar{z}$. If $\|d \log \lambda\|_{\omega_{D}}^{2} \equiv 4$, then $D$ is affine equivalent to the half-plane $\mathbf{H}=\{z \in \mathbb{C}: \operatorname{Re} z<0\}$.

## 3. Existence of nowhere vanishing complete holomorphic vector field

In this section, we will study an existence of a complete holomorphic tangent vector field on a Riemann surface $X$ which admits a Kähler potential of the Poincaré metric with a constant length differential.

By a holomorphic tangent vector field of a Riemann surface $X$, we means a holomorphic section $\mathcal{W}$ to the holomorphic tangent bundle $T^{1,0} X$. If the corresponding real tangent vector field $\operatorname{Re} \mathcal{W}=\mathcal{W}+\overline{\mathcal{W}}$ is complete, we also say $\mathcal{W}$ is complete. Thus the complete holomorphic tangent vector field generates a 1-parameter family of holomorphic transformations.

In this section, we will show that:
Theorem 3.1. Let $X$ be a Riemann surface with the Poincaré metric $\omega_{X}$. If there is a function $\varphi: X \rightarrow \mathbb{R}$ with

$$
\begin{equation*}
d d^{c} \log \varphi=2 \omega_{X} \quad \text { and } \quad\|d \log \varphi\|_{\omega_{X}}^{2} \equiv 4 \tag{3.1}
\end{equation*}
$$

then there is a nowhere vanishing complete holomorphic vector field $\mathcal{W}$ such that $(\operatorname{Re} \mathcal{W}) \varphi \equiv 0$.

Proof. Take a local holomorphic coordinate function $z$ and let $\omega_{X}=i \lambda d z \wedge d \bar{z}$. The equation (3.1) can be written by

$$
(\log \varphi)_{z \bar{z}}=2 \lambda \quad \text { and } \quad(\log \varphi)_{z}(\log \varphi)_{\bar{z}}=4 \lambda
$$

Here, $(\log \varphi)_{z}=\frac{\partial}{\partial z} \log \varphi,(\log \varphi)_{\bar{z}}=\frac{\partial}{\partial \bar{z}} \log \varphi$ and $(\log \varphi)_{z \bar{z}}=\frac{\partial^{2}}{\partial z \partial \bar{z}} \log \varphi$. This implies that

$$
\begin{aligned}
\left(\varphi^{-1 / 2}\right)_{z} & =\frac{\partial}{\partial z} \varphi^{-1 / 2}=-\frac{1}{2} \varphi^{-1 / 2}(\log \varphi)_{z} \\
\left(\varphi^{-1 / 2}\right)_{z \bar{z}} & =\frac{\partial^{2}}{\partial z \partial \bar{z}} \varphi^{-1 / 2}=-\frac{1}{2} \varphi^{-1 / 2}(\log \varphi)_{z \bar{z}}+\frac{1}{4} \varphi^{-1 / 2}(\log \varphi)_{z}(\log \varphi)_{\bar{z}} \\
& =-\frac{1}{2} \varphi^{-1 / 2}\left((\log \varphi)_{z \bar{z}}-\frac{1}{2}(\log \varphi)_{z}(\log \varphi)_{\bar{z}}\right) \\
& =0 .
\end{aligned}
$$

Thus we have that the function $\varphi^{-1 / 2}$ is harmonic so $\left(\varphi^{-1 / 2}\right)_{z}$ is holomorphic.
Let us consider a local holomorphic vector field,

$$
\mathcal{W}=\frac{i}{\left(\varphi^{-1 / 2}\right)_{z}} \frac{\partial}{\partial z}=\frac{-2 i \varphi^{3 / 2}}{\varphi_{z}} \frac{\partial}{\partial z}=\frac{-2 i \varphi^{1 / 2}}{(\log \varphi)_{z}} \frac{\partial}{\partial z}
$$

In any other local holomorphic coordinate function $w$, we have

$$
\mathcal{W}=\frac{i}{\left(\varphi^{-1 / 2}\right)_{z}} \frac{\partial}{\partial z}=\frac{i}{\left(\varphi^{-1 / 2}\right)_{w} \frac{\partial w}{\partial z}} \frac{\partial w}{\partial z} \frac{\partial}{\partial w}=\frac{i}{\left(\varphi^{-1 / 2}\right)_{w}} \frac{\partial}{\partial w},
$$

so $W$ is globally defined on $X$. Now we will show that $\mathcal{W}$ satisfies conditions in the theorem.

Since

$$
\left\|\varphi^{-1 / 2} \mathcal{W}\right\|_{\omega_{X}}^{2}=\left\|\frac{-2 i}{(\log \varphi)_{z}} \frac{\partial}{\partial z}\right\|_{\omega_{X}}^{2}=\frac{4 \lambda}{(\log \varphi)_{z}(\log \varphi)_{\bar{z}}}=1
$$

the vector field $\varphi^{-1 / 2} \mathcal{W}$ has a unit length with respect to the complete metric $\omega_{X}$, so the corresponding real vector field $\operatorname{Re} \varphi^{-1 / 2} \mathcal{W}=\varphi^{-1 / 2}(\mathcal{W}+\overline{\mathcal{W}})$ is complete. Moreover

$$
(\operatorname{Re} \mathcal{W}) \varphi=\frac{-2 i \varphi^{3 / 2}}{\varphi_{z}} \varphi_{z}+\frac{2 i \varphi^{3 / 2}}{\varphi_{\bar{z}}} \varphi_{\bar{z}}=0
$$

Hence it remains to show the completeness of $\mathcal{W}$. Take any integral curve $\gamma: \mathbb{R} \rightarrow X$ of $\varphi^{-1 / 2} \operatorname{Re} \mathcal{W}$. It satisfies

$$
\left(\varphi^{-1 / 2}(\operatorname{Re} \mathcal{W})\right) \circ \gamma=\dot{\gamma}
$$

equivalently

$$
(\operatorname{Re} \mathcal{W}) \circ \gamma=\left(\varphi^{1 / 2} \circ \gamma\right) \dot{\gamma}
$$

The condition $(\operatorname{Re} \mathcal{W}) \varphi \equiv 0$, equivalently $\varphi^{-1 / 2}(\operatorname{Re} \mathcal{W}) \varphi \equiv 0$, implies that the curve $\gamma$ is on a level set of $\varphi$ so $\varphi^{1 / 2} \circ \gamma \equiv C$ for some constant $C$. The curve $\sigma: \mathbb{R} \rightarrow X$ defined by $\sigma(t)=\gamma(C t)$ satisfies

$$
(\operatorname{Re} \mathcal{W}) \circ \sigma(t)=(\operatorname{Re} \mathcal{W})(\gamma(C t))=C \dot{\gamma}(C t)=\dot{\sigma}(t)
$$

This means that $\sigma: \mathbb{R} \rightarrow X$ is the integral curve of $\operatorname{Re} \mathcal{W}$; therefore $\operatorname{Re} \mathcal{W}$ is complete. This completes the proof.

## 4. Complete holomorphic vector fields on the unit disc

In this section, we introduce parabolic and hyperbolic vector fields on the unit disc and discuss their relation to the model potential,

$$
\begin{equation*}
\varphi_{0}=\frac{|1+z|^{4}}{\left(1-|z|^{2}\right)^{2}} \tag{4.1}
\end{equation*}
$$

where it is $\varphi_{\theta}$ in (2.1) with $\theta=0$.

### 4.1. Nowhere vanishing complete holomorphic vector fields from the left-half plane

On the left-half plane $\mathbf{H}=\{w \in \mathbb{C}: \operatorname{Re} w<0\}$, there are two kinds of affine transformations:

$$
\mathcal{D}_{s}(w)=e^{2 s} w \quad \text { and } \quad \mathcal{T}_{s}(w)=w+2 i s
$$

for $s \in \mathbb{R}$. Their infinitesimal generators are

$$
\mathcal{D}=2 w \frac{\partial}{\partial w} \quad \text { and } \quad \mathcal{T}=2 i \frac{\partial}{\partial w}
$$

which are nowhere vanishing complete holomorphic vector fields of $\mathbf{H}$. Note that

$$
\begin{equation*}
\left(\mathcal{T}_{s}\right)_{*} \mathcal{D}=2(w-2 i s) \frac{\partial}{\partial w}=\mathcal{D}-2 s \mathcal{T} \quad \text { and } \quad\left(\mathcal{T}_{s}\right)_{*} \mathcal{T}=2 i \frac{\partial}{\partial w}=\mathcal{T} \tag{4.2}
\end{equation*}
$$

for any $s$.
For the Cayley transform $F: \mathbf{H} \rightarrow \Delta$ defined by

$$
\begin{align*}
& F: \mathbf{H} \longrightarrow \Delta \\
& w \longmapsto z=\frac{1+w}{1-w} \tag{4.3}
\end{align*}
$$

we can take two nowhere vanishing complete holomorphic vector fields of $\Delta$ :

$$
\mathcal{H}=F_{*}(\mathcal{D})=\left(z^{2}-1\right) \frac{\partial}{\partial z}
$$

and

$$
\mathcal{P}=F_{*}(\mathcal{T})=i(z+1)^{2} \frac{\partial}{\partial z}
$$

When we define $\mathcal{H}_{s}=F \circ \mathcal{D}_{s} \circ F^{-1}$ and $\mathcal{P}_{s}=F \circ \mathcal{T}_{s} \circ F^{-1}$, vector fields $\mathcal{H}$ and $\mathcal{P}$ are infinitesimal generators of $\mathcal{H}_{s}$ and $\mathcal{P}_{s}$, respectively. Moreover Equation (4.2) can be written by

$$
\begin{equation*}
\left(\mathcal{P}_{s}\right)_{*} \mathcal{H}=\mathcal{H}-2 s \mathcal{P} \quad \text { and } \quad\left(\mathcal{P}_{s}\right)_{*} \mathcal{P}=\mathcal{P} . \tag{4.4}
\end{equation*}
$$

There is another complete holomorphic vector field $\mathcal{R}=i z \partial / \partial z$ generating the rotational symmetry

$$
\begin{equation*}
\mathcal{R}_{s}(z)=e^{i s} z \tag{4.5}
\end{equation*}
$$

The holomorphic automorphism group of $\Delta$ is a real 3-dimension connected Lie group (cf. see $[1,9]$ ), we can conclude that any complete holomorphic vector field can be a real linear combination of $\mathcal{H}, \mathcal{P}$ and $\mathcal{R}$. Since $\mathcal{H}(-1)=\mathcal{P}(-1)=0$ and $\mathcal{R}(-1)=-i \partial / \partial z$, we have:
Lemma 4.1. If $\mathcal{W}$ is a complete holomorphic vector field of $\Delta$ satisfying $\mathcal{W}(-1)=0$, then there exist $a, b \in \mathbb{R}$ with $\mathcal{W}=a \mathcal{H}+b \mathcal{P}$.

### 4.2. Hyperbolic vector fields

In this subsection, we will show that the hyperbolic vector field $\mathcal{H}$ can not be tangent to a Kähler potential with a constant length differential.

By the simple computation,

$$
\mathcal{H}\left(\log \varphi_{0}\right)=\left(z^{2}-1\right) \frac{2(1+\bar{z})}{(1+z)\left(1-|z|^{2}\right)}=2 \frac{|z|^{2}+z-\bar{z}-1}{\left(1-|z|^{2}\right)}
$$

we get

$$
(\operatorname{Re} \mathcal{H}) \log \varphi_{0} \equiv-4
$$

That means $\operatorname{Re} \mathcal{H}$ is nowhere tangent to $\varphi_{0}$. Moreover, we have:
Lemma 4.2. Let $\varphi: \Delta \rightarrow \mathbb{R}$ with $d d^{c} \log \varphi=2 \omega_{\Delta}$ and $\|d \log \varphi\|_{\omega_{\Delta}}^{2} \equiv 4$. If $(\operatorname{Re} \mathcal{H}) \log \varphi \equiv c$ for some $c$, then $c= \pm 4$.
Proof. Since $d d^{c} \log \varphi_{0}=2 \omega_{\Delta}$ also, the function $\log \varphi-\log \varphi_{0}$ is harmonic; hence we may let $\log \varphi=\log \varphi_{0}+f+\bar{f}$ for some holomorphic function $f: \Delta \rightarrow$ $\mathbb{C}$. Then the condition $(\operatorname{Re} \mathcal{H}) \log \varphi \equiv c$ can be written by

$$
\begin{equation*}
(\operatorname{Re} \mathcal{H}) \log \varphi=-4+\left(z^{2}-1\right) f^{\prime}+\left(\bar{z}^{2}-1\right) \bar{f}^{\prime} \equiv c \tag{4.6}
\end{equation*}
$$

This implies that $\left(z^{2}-1\right) f^{\prime}$ is constant. Thus we can let

$$
\begin{equation*}
f^{\prime}=\frac{C}{z^{2}-1} \tag{4.7}
\end{equation*}
$$

for some $C \in \mathbb{C}$. Since

$$
\frac{\partial}{\partial z} \log \varphi=f^{\prime}+\frac{\partial}{\partial z} \log \varphi_{0}=f^{\prime}+\frac{2(1+\bar{z})}{(1+z)\left(1-|z|^{2}\right)}
$$

we have

$$
\|d \log \varphi\|_{\omega_{\Delta}}^{2}=\left(\frac{\partial}{\partial z} \log \varphi\right)\left(\frac{\partial}{\partial \bar{z}} \log \varphi\right) \frac{1}{\lambda_{\Delta}}
$$

$=\left|f^{\prime}\right|^{2}\left(1-|z|^{2}\right)^{2}+\frac{2(1+\bar{z})\left(1-|z|^{2}\right)}{(1+z)} \bar{f}^{\prime}+\frac{2(1+z)\left(1-|z|^{2}\right)}{(1+\bar{z})} f^{\prime}+\left\|d \log \varphi_{0}\right\|_{\omega_{\Delta}}^{2}$.
From the condition $\|d \log \varphi\|_{\omega_{\Delta}}^{2} \equiv 4 \equiv\left\|d \log \varphi_{0}\right\|_{\omega_{\Delta}}^{2}$, it follows

$$
\left|f^{\prime}\right|^{2}\left(1-|z|^{2}\right)^{2}=-\frac{2(1+\bar{z})\left(1-|z|^{2}\right)}{(1+z)} \bar{f}^{\prime}-\frac{2(1+z)\left(1-|z|^{2}\right)}{(1+\bar{z})} f^{\prime},
$$

equivalently

$$
\begin{equation*}
\frac{1}{2}\left|f^{\prime}\right|^{2}\left(1-|z|^{2}\right)=-\frac{(1+\bar{z})}{(1+z)} \bar{f}^{\prime}-\frac{(1+z)}{(1+\bar{z})} f^{\prime} \tag{4.8}
\end{equation*}
$$

Applying (4.7) to the right side above,

$$
\begin{aligned}
-\frac{(1+\bar{z})}{(1+z)} \bar{f}^{\prime}-\frac{(1+z)}{(1+\bar{z})} f^{\prime} & =\frac{(1+\bar{z})}{(1+z)} \frac{\bar{C}}{1-\bar{z}^{2}}+\frac{(1+z)}{(1+\bar{z})} \frac{C}{1-z^{2}} \\
& =\frac{\left(1+\bar{z}-z-|z|^{2}\right) \bar{C}+\left(1-\bar{z}+z-|z|^{2}\right) C}{\left|1-z^{2}\right|^{2}} .
\end{aligned}
$$

Let $C=a+b i$ for $a, b \in \mathbb{R}$, then

$$
\left(1+\bar{z}-z-|z|^{2}\right) \bar{C}+\left(1-\bar{z}+z-|z|^{2}\right) C=2 a\left(1-|z|^{2}\right)+2 b i(z-\bar{z})
$$

Now Equation (4.8) can be written by

$$
\frac{1}{2} \frac{|C|^{2}}{\left|z^{2}-1\right|^{2}}\left(1-|z|^{2}\right)=\frac{2 a\left(1-|z|^{2}\right)+2 b i(z-\bar{z})}{\left|1-z^{2}\right|^{2}}
$$

so we have

$$
\left(|C|^{2}-4 a\right)\left(1-|z|^{2}\right)=4 b i(z-\bar{z})
$$

on $\Delta$. Take $\partial \bar{\partial}$ to above, we have

$$
|C|^{2}-4 a=0
$$

Simultaneously $b=0$ so $C=a$. Now we have $a^{2}=4 a$. Such $a$ is 0 or 4. If $f^{\prime}=4 /\left(z^{2}-1\right)$, then $c=4$ from (4.6). If $f^{\prime}=0$, then $c=-4$.

### 4.3. Parabolic vector fields

Since

$$
\mathcal{P}\left(\log \varphi_{0}\right)=i(z+1)^{2} \frac{2(1+\bar{z})}{(1+z)\left(1-|z|^{2}\right)}=2 i \frac{|1+z|^{2}}{1-|z|^{2}}
$$

we have

$$
(\operatorname{Re} \mathcal{P}) \log \varphi_{0} \equiv 0
$$

That means that the parabolic vector field $\mathcal{P}$ is tangent to $\varphi_{0}$. The vector field $\mathcal{P}$ is indeed the nowhere vanishing complete holomorphic vector field as constructed in Theorem 3.1 corresponding to $\varphi_{0}$. The main result of this section is the following.

Lemma 4.3. Let $\varphi: \Delta \rightarrow \mathbb{R}$ with $d d^{c} \log \varphi=2 \omega_{\Delta}$ and $\|d \log \varphi\|_{\omega_{\Delta}}^{2} \equiv 4$. If $(\operatorname{Re} \mathcal{P}) \log \varphi \equiv c$ for some $c$, then $c=0$ and $\varphi=r \varphi_{0}$ for some $r>0$.

Proof. By the same way in the proof of Lemma 4.2, we let $\log \varphi=\log \varphi_{0}+f+\bar{f}$ for some holomorphic $f: \Delta \rightarrow \mathbb{C}$. Since

$$
\begin{equation*}
(\operatorname{Re} \mathcal{P}) \log \varphi=i(z+1)^{2} f^{\prime}-i(\bar{z}+1)^{2} \bar{f}^{\prime} \equiv c \tag{4.9}
\end{equation*}
$$

it follows that $(z+1)^{2} f^{\prime}$ is constant. Thus we have

$$
\begin{equation*}
f^{\prime}=\frac{C}{(z+1)^{2}} \tag{4.10}
\end{equation*}
$$

for some $C \in \mathbb{C}$. Since (4.8) also holds, we can apply (4.10) to the right side of (4.8) to get

$$
\begin{aligned}
-\frac{(1+\bar{z})}{(1+z)} \bar{f}^{\prime}-\frac{(1+z)}{(1+\bar{z})} f^{\prime} & =-\frac{(1+\bar{z})}{(1+z)} \frac{\bar{C}}{(\bar{z}+1)^{2}}-\frac{(1+z)}{(1+\bar{z})} \frac{C}{(z+1)^{2}} \\
& =\frac{-\bar{C}}{|1+z|^{2}}+\frac{-C}{|1+z|^{2}}=\frac{-\bar{C}-C}{|1+z|^{2}}
\end{aligned}
$$

Now Equation (4.8) is can be written by

$$
\frac{|C|^{2}}{|z+1|^{4}}\left(1-|z|^{2}\right)=2 \frac{-\bar{C}-C}{|1+z|^{2}}
$$

equivalently

$$
|C|^{2}\left(1-|z|^{2}\right)=-(2 \bar{C}+2 C)|1+z|^{2}
$$

Evaluating $z=0$, we have $|C|^{2}=-2 \bar{C}-2 C$. And taking $\partial \bar{\partial}$ to above, we have $-|C|^{2}=-2 \bar{C}-2 C$. It follows that $C=0$ so $f$ is constant. Moreover Equation (4.9) implies that $c=0$.

## 5. Proof of the main theorem

Now we prove Theorem 2.2 and Corollary 2.3.
Proof of Theorem 2.2. Let $\varphi: \Delta \rightarrow \mathbb{R}$ be a function with

$$
d d^{c} \log \varphi=2 \omega_{\Delta} \quad \text { and } \quad\|d \log \varphi\|_{\omega_{\Delta}}^{2} \equiv 4
$$

By Theorem 3.1, we can take a nowhere vanishing complete holomorphic vector field $\mathcal{W}$ with $(\operatorname{Re} \mathcal{W}) \varphi \equiv 0$. Since every automorphism of $\Delta$ has at least one fixed point on $\bar{\Delta}$ and $\mathcal{W}$ is nowhere vanishing on $\Delta$, any nontrivial automorphism generated by $\operatorname{Re} \mathcal{W}$ has no fixed point in $\Delta$ and should have a common fixed point $p$ at the boundary $\partial \Delta$. This means $p$ is a vanishing point of $\mathcal{W}$. Consider a rotational symmetry $\mathcal{R}_{\theta}$ in (4.5) satisfying $\mathcal{R}_{\theta}(-1)=p$. We will show that $\varphi \circ \mathcal{R}_{\theta}=r \varphi_{0}$ where $\varphi_{0}$ is as in (4.1) and $r>0$. This implies that $\varphi=r \varphi_{-\theta}$.

Now we can simply denote by $\varphi=\varphi \circ \mathcal{R}_{\theta}$ and $\mathcal{W}=\left(\mathcal{R}_{\theta}^{-1}\right)_{*} \mathcal{W}$. Since -1 is a vanishing point of $\mathcal{W}$, Lemma 4.1 implies

$$
\mathcal{W}=a \mathcal{H}+b \mathcal{P}
$$

for some real numbers $a, b$.

Suppose that $a \neq 0$. Equation (4.4) implies that

$$
\left(\mathcal{P}_{s}\right)_{*} \mathcal{W}=\left(\mathcal{P}_{s}\right)_{*}(a \mathcal{H}+b \mathcal{P})=a \mathcal{H}-2 a s \mathcal{P}+b \mathcal{P}=a \mathcal{H}+(b-2 a s) \mathcal{P}
$$

Take $s=b / 2 a$, then $\widetilde{\mathcal{W}}=\left(\mathcal{P}_{s}\right)_{*} \mathcal{W}=a \mathcal{H}$. Let $\tilde{\varphi}=\varphi \circ \mathcal{P}_{-s}$ for this $s$. Then $\tilde{\varphi}$ satisfies conditions in Theorem 2.2 and $(\operatorname{Re} \widetilde{\mathcal{W}}) \tilde{\varphi} \equiv 0$. But Lemma 4.2 said that $(\operatorname{Re} \widetilde{\mathcal{W}}) \tilde{\varphi}=a(\operatorname{Re} \mathcal{H}) \tilde{\varphi} \equiv \pm 4 a \tilde{\varphi}$. It contradicts to $(\operatorname{Re} \mathcal{W}) \varphi \equiv 0$ equivalently $(\operatorname{Re} \widetilde{\mathcal{W}}) \tilde{\varphi} \equiv 0$. Thus $a=0$.

Now $\mathcal{W}=b \mathcal{P}$. Since $\mathcal{W}$ is nowhere vanishing already, $b \neq 0$. The condition $(\operatorname{Re} \mathcal{W}) \varphi \equiv 0$ implies $(\operatorname{Re} \mathcal{P}) \varphi \equiv 0$. Lemma 4.3 says that $\varphi=r \varphi_{0}$ for some positive $r$. This completes the proof.

Proof of Corollary 2.3. Let $D$ be a simply connected proper domain in $\mathbb{C}$ and let $\omega_{D}=i \lambda_{D} d z \wedge d \bar{z}$ be its Poincaré metric with $\left\|d \log \lambda_{D}\right\|_{\omega_{D}}^{2} \equiv 4$. By Theorem 3.1, there is a nowhere vanishing complete holomorphic vector field $\mathcal{W}$ with $(\operatorname{Re} \mathcal{W}) \lambda_{D} \equiv 0$. Take a biholomorphism $G: \Delta \rightarrow D$ and let

$$
\varphi=\lambda_{D} \circ G \quad \text { and } \quad \mathcal{Z}=\left(G^{-1}\right)_{*} \mathcal{W}
$$

Note that $(\operatorname{Re} \mathcal{Z}) \varphi \equiv 0$ by assumption. Using the rotational symmetry $\mathcal{R}_{\theta}$ of $\Delta$ which is also affine, we may assume that $\mathcal{Z}(-1)=0$ and we will prove that $G$ is a Cayley transform.

Since $G:\left(\Delta, \omega_{\Delta}\right) \rightarrow\left(D, \omega_{D}\right)$ is an isometry, we have $G^{*} \omega_{D}=\omega_{\Delta}$, equivalently

$$
\varphi=\frac{\lambda_{\Delta}}{\left|G^{\prime}\right|^{2}}
$$

Moreover $d \log \varphi=d\left(G^{*} \log \lambda_{D}\right)$ implies that $\|d \log \varphi\|_{\omega_{D}}^{2}=\left\|d\left(G^{*} \log \lambda_{D}\right)\right\|_{\omega_{D}}^{2}$ $\equiv 4$. By Theorem 2.2, we have

$$
\frac{\lambda_{\Delta}}{\left|G^{\prime}\right|^{2}}=\varphi=r \varphi_{0}=r \lambda_{\Delta}|1+z|^{4}
$$

for some positive $r$. This means that $G^{\prime}=e^{i \theta^{\prime}} / \sqrt{r}(1+z)^{2}$ for some $\theta^{\prime} \in \mathbb{R}$ so that

$$
G=\frac{e^{i \theta^{\prime}}}{2 \sqrt{r}} \frac{z-1}{z+1}+C .
$$

Since the function $z \mapsto(z-1) /(z+1)$ is the inverse mapping of the Cayley transform $F: \mathbf{H} \rightarrow \Delta$ in (4.3), we have

$$
\begin{aligned}
G \circ F: \mathbf{H} & \rightarrow D \\
& z \mapsto \frac{e^{i \theta^{\prime}}}{2 \sqrt{r}} z+C .
\end{aligned}
$$

This implies that $D=G(F(\mathbf{H}))$ is affine equivalent to $\mathbf{H}$.

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